

Frequency Corrections in Flux Computations



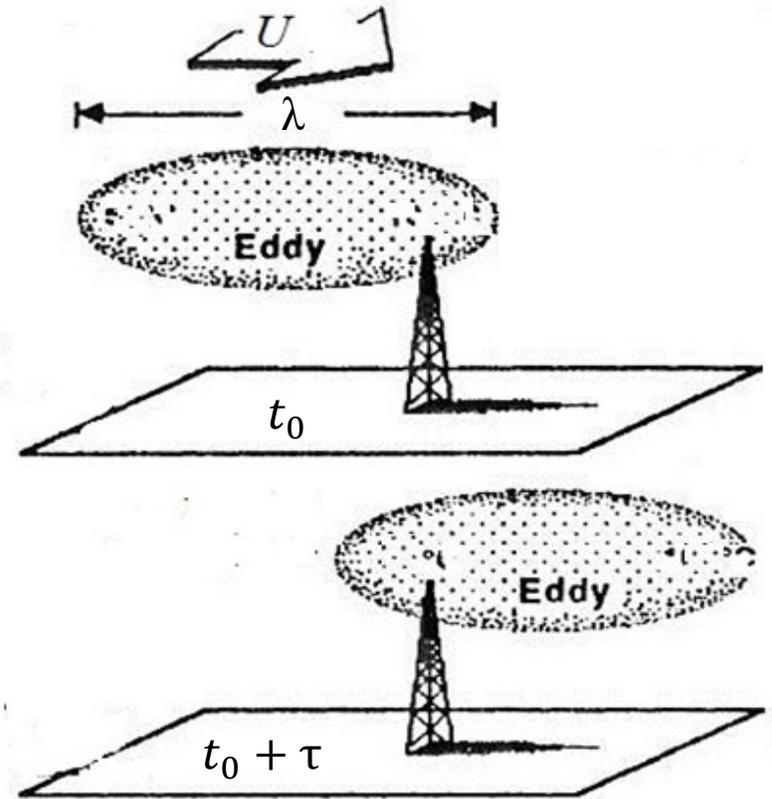
Xinhua Zhou, Ph.D., Senior Scientist
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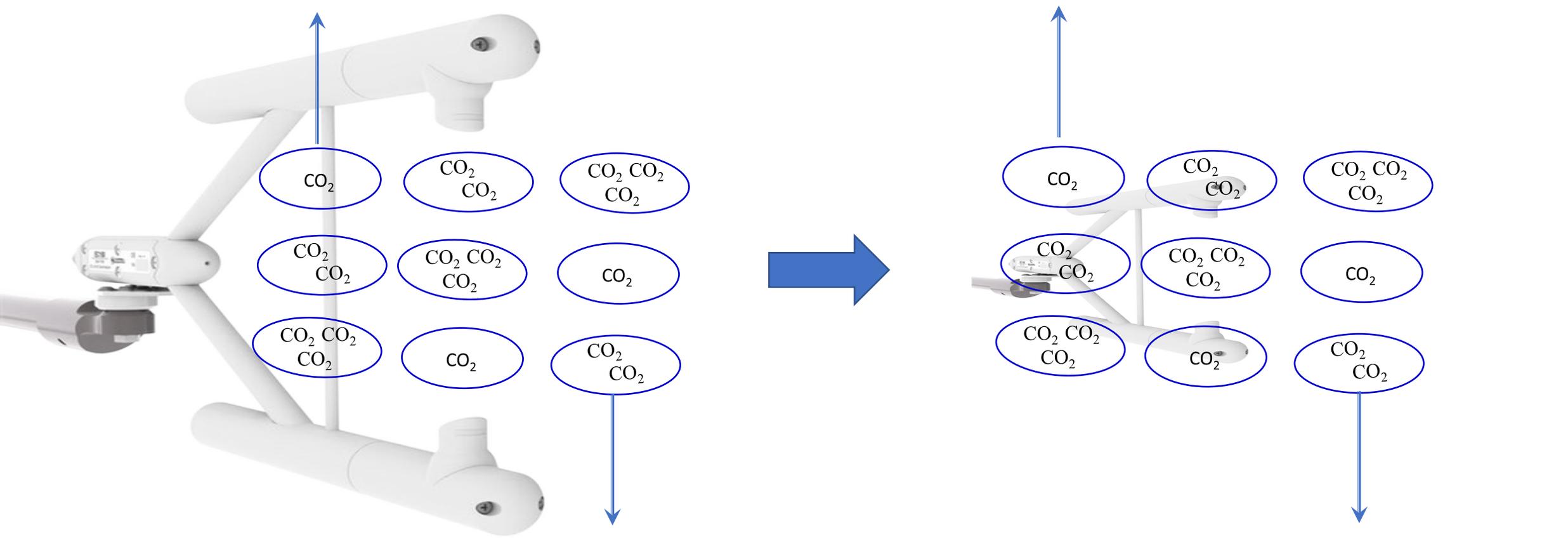
Eddy sizes, speed, and frequency

$$\text{Period}(\tau) = \frac{\text{Wavelength}(\lambda)}{\text{Wind speed}(U)} = \frac{1}{\text{Frequency}(f)}$$

$$f = \frac{U}{\lambda}$$



Measuring CO₂ for the same air flow by two gas analyzers with different path lengths



| | | | |
|---|---------------|-----------|---------------|
| CO ₂ (ρ_{CO_2} in molCO ₂ m ⁻³) | 2 | 2 | 2 |
| Vertical wind (w in m s ⁻¹) | $\bar{w} + 1$ | \bar{w} | $\bar{w} - 1$ |
| $\rho_{CO_2} w$ (molCO ₂ m ⁻² s ⁻¹) | 0 | 0 | 0 |

$$F_c = \overline{\rho_{CO_2} w'}^a$$

| | | | |
|---|---------------|-----------|---------------|
| CO ₂ (ρ_{CO_2} in molCO ₂ m ⁻³) | 2 | 3 | 1 |
| Vertical wind (w in m s ⁻¹) | $\bar{w} + 1$ | \bar{w} | $\bar{w} - 1$ |
| $\rho_{CO_2} w$ (molCO ₂ m ⁻² s ⁻¹) | 0 | 0 | +1 |

$$F_c = \overline{\rho_{CO_2} w'}^a$$

^a Air density is assumed to be a constant.



551.524.4 : 551.551.8 : 517.512.2

The effect of vertical line averaging on the spectra of
temperature and heat-flux

By J. C. KAIMAL .

Air Force Cambridge Research Laboratories, Bedford, Massachusetts

1968

Quart. J. Roy. Meteor
Soc. 94: 149-155

(Manuscript received 30 June 1967; in revised form 13 November 1967)

SUMMARY

The paper describes an attempt to evaluate the effects of vertical line averaging on temperature statistics by using simultaneous measurements from sensors of different lengths. Results show that the transfer function for vertical line averaging derived by Gurvich for the vertical velocity component applies also to temperature spectra. Heat-flux, when computed with vertical velocity measurements from a sonic anemometer, is relatively insensitive to line averaging (at least up to 80 cm at 9 m height). But the coherence between the line average and its mid-point value shows consistent behaviour only for short line lengths of the order of 20 cm. With longer lengths the coherence is likely to vary from one run to another depending on mean wind speed.

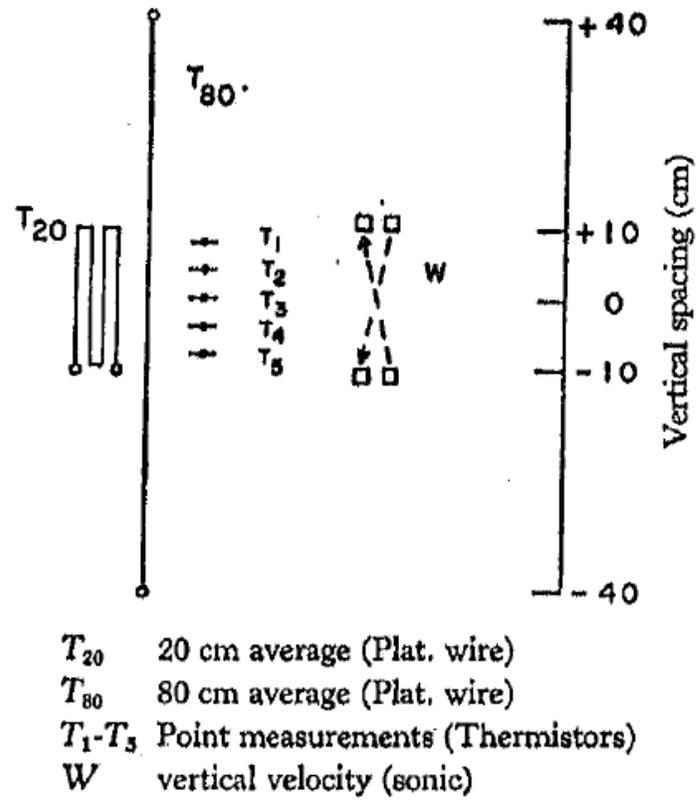


Figure 1. Diagram of sensor placement for the experiment.

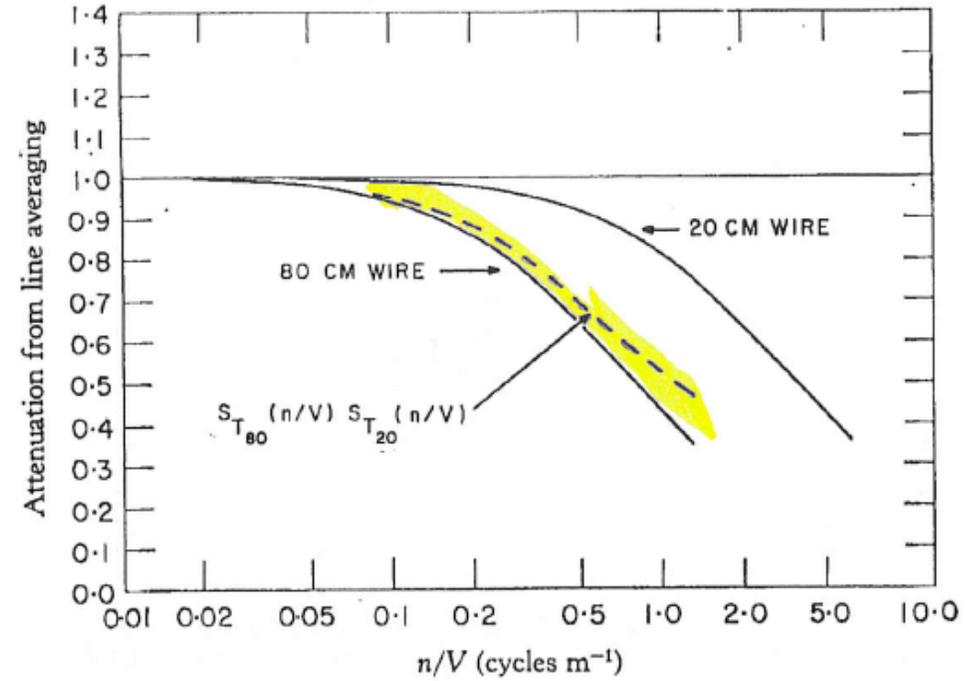


Figure 3. Theoretical attenuation curves for 20- and 80-cm line averaging. Relative attenuation between the two is indicated by dashed curves.

Implication of time constant (τ) in physics

$$\tau = \gamma D^2 \frac{\rho_a C}{k_a \text{Nu}}$$

γ Shape factor

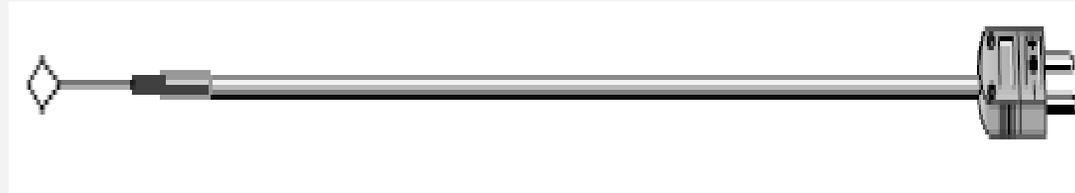
D Dimension

ρ_a Material density of sensor sensing component

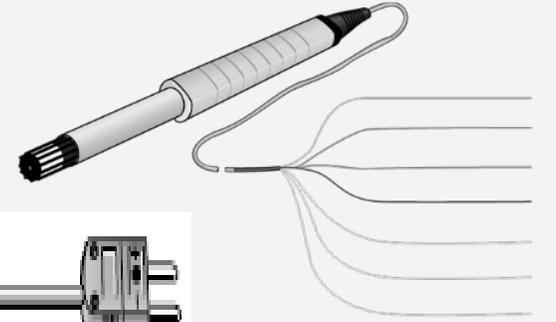
C Specific heat of materials for sensor sensing

k_a Thermal conductivity of air

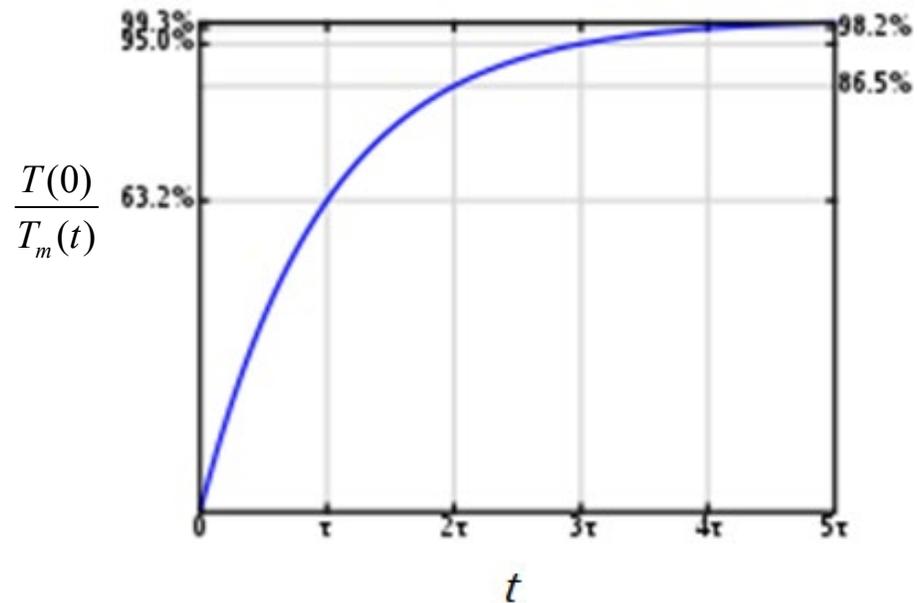
Nu Nusselt number



Moore (1986)



Response of temperature sensor as described in terms of time constant



$$\frac{dT_m(t)}{dt} = \frac{T(t) - T_m(t)}{\tau}$$

t Time

$T(t)$ Temperature of media at t

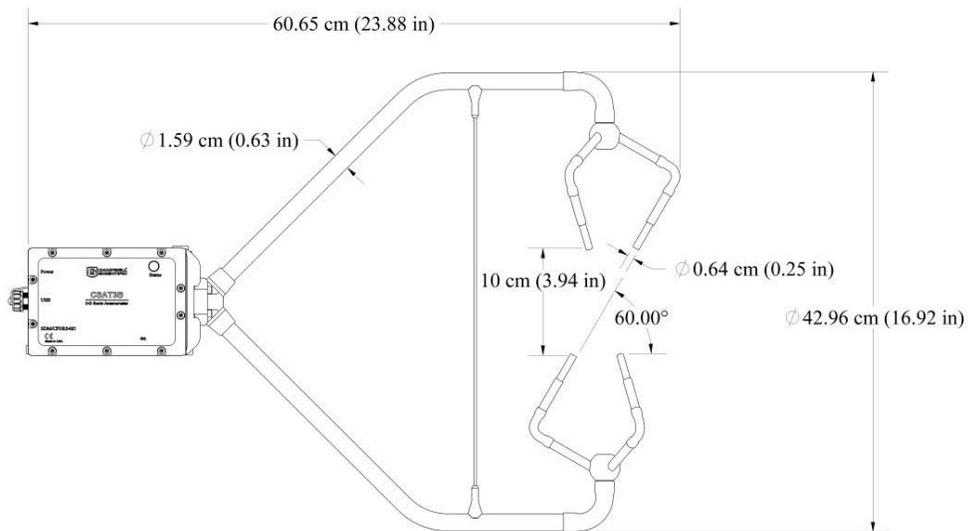
$T_m(t)$ Temperature measured at t

τ Time constant

Burt and Podesta (2020)

Response time of sensors

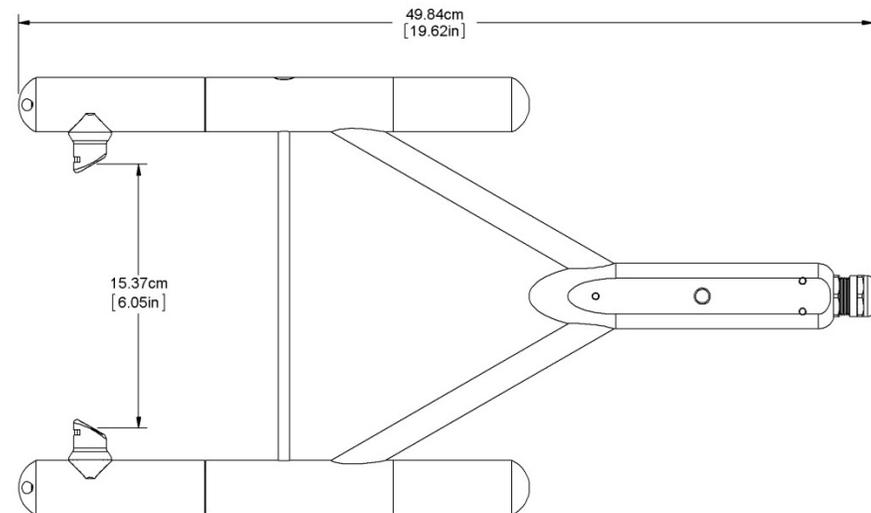
Campbell Scientific sonic anemometer thermometry
(CSAT, Campbell Scientific Inc., UT, USA)



$$\text{Response time}^a = \frac{\text{Path length (11.54 cm)}}{\text{Speed of sound} - \text{flow speed along the path}}$$
$$\approx 3.6 \times 10^{-4} \text{ s}$$

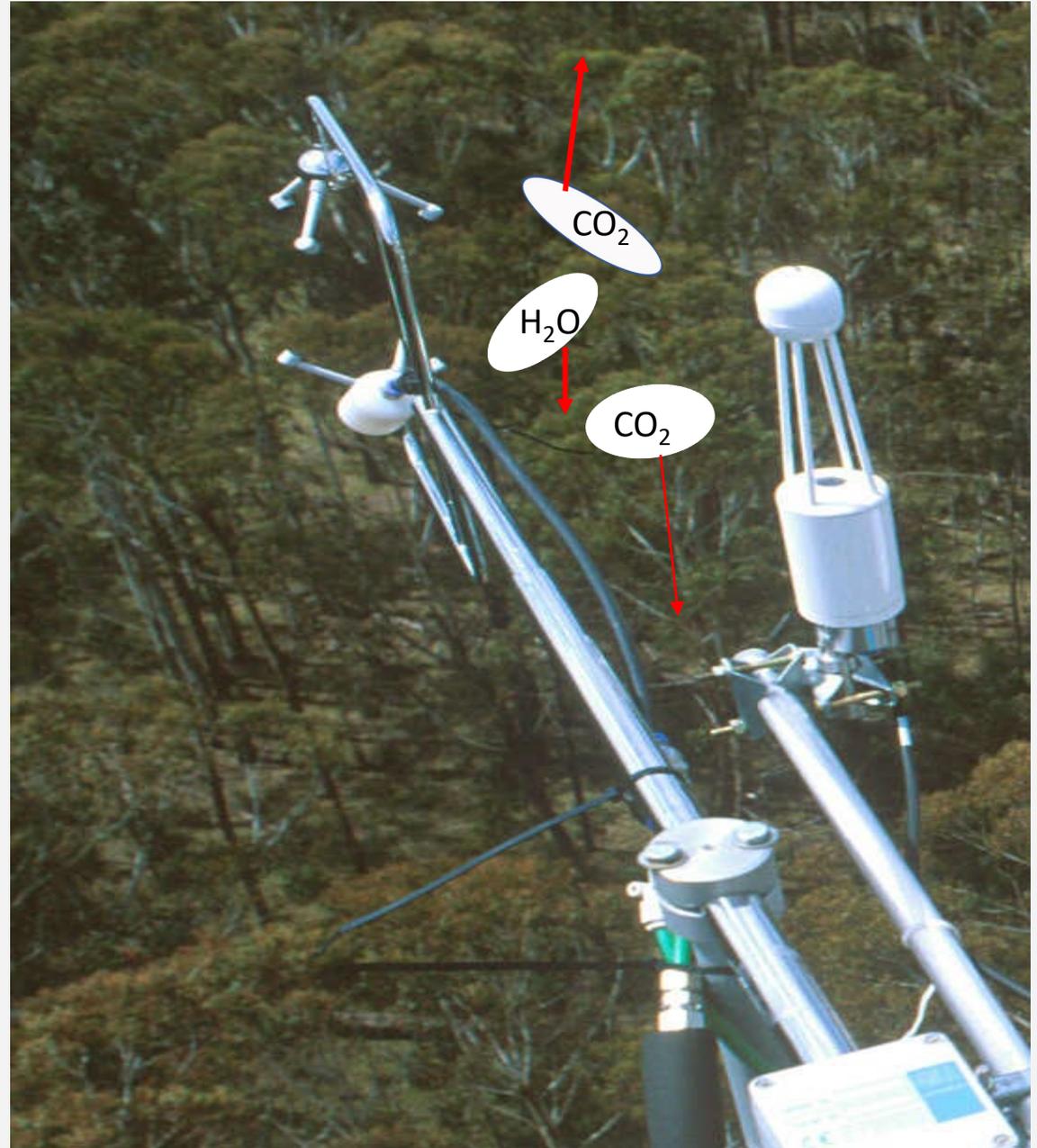
^a Electronic response is assumed to be instantaneous.

Infrared CO₂-H₂O analyzer
(EC150, Campbell Scientific Inc., UT, USA)



$$\text{Response time}^a = \frac{\text{Path length}}{\text{Speed of light}}$$
$$\approx 5.1 \times 10^{-10} \text{ s}$$

Frequency loss due to sensor separation



Low frequency loss due to block averaging in flux computations

$$\text{CO}_2 \quad F_{\text{CO}_2} = \rho_d \overline{w' \chi'_{\text{CO}_2}} \quad (\text{e.g., } w' = w - \bar{w})$$

$$\text{H}_2\text{O} \quad F_{\text{H}_2\text{O}} = \rho_d \overline{w' \chi'_{\text{H}_2\text{O}}}$$

Summary for frequency losses

High frequency

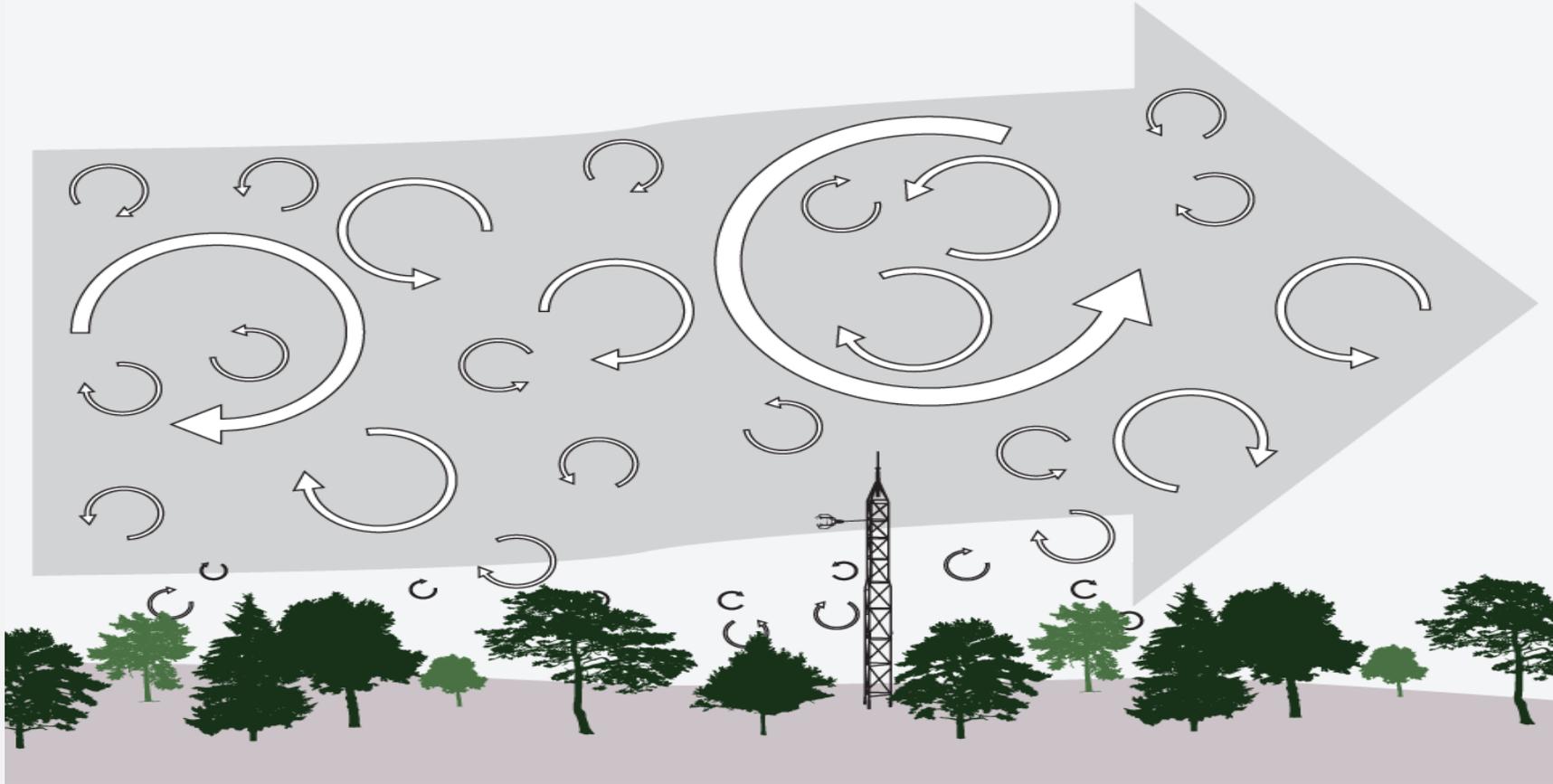
- a. Line averaging.
- b. Delay in response to measured variables.
- c. Spatial separation between sensors

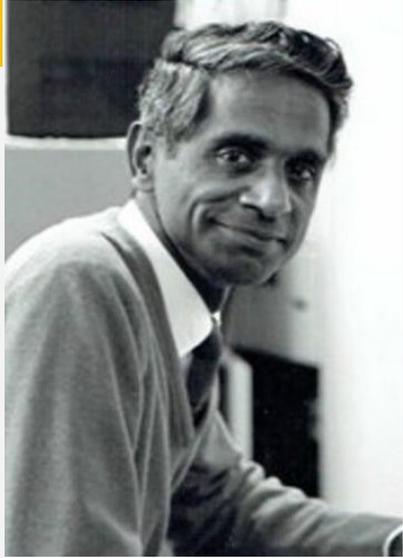
Low frequency

- a. Block averaging.

Any variable measured in an eddy-covariance system is a time series (i.e., a function of time).

For frequency corrections, this time series in a time domain needs to be transformed into a frequency domain.





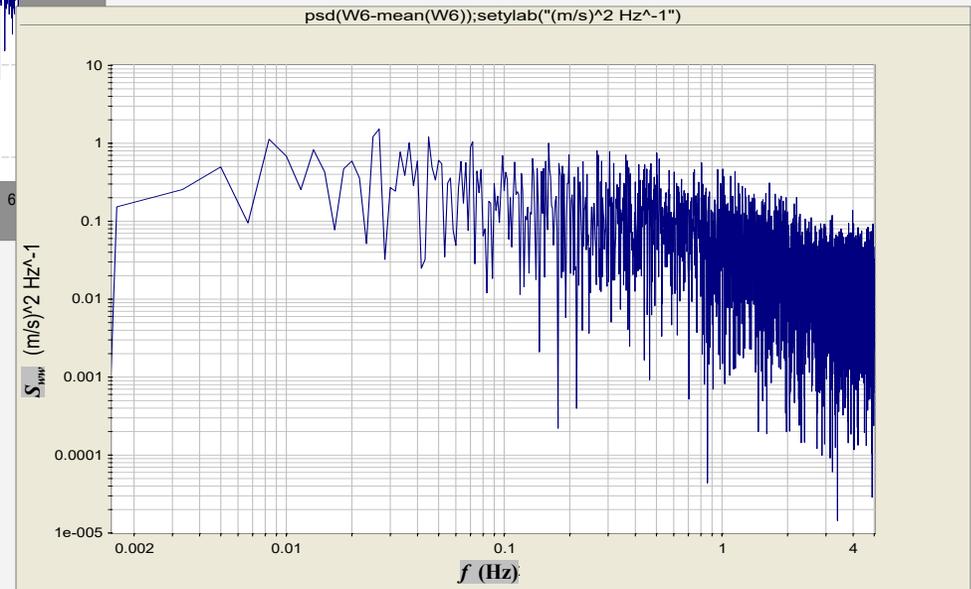
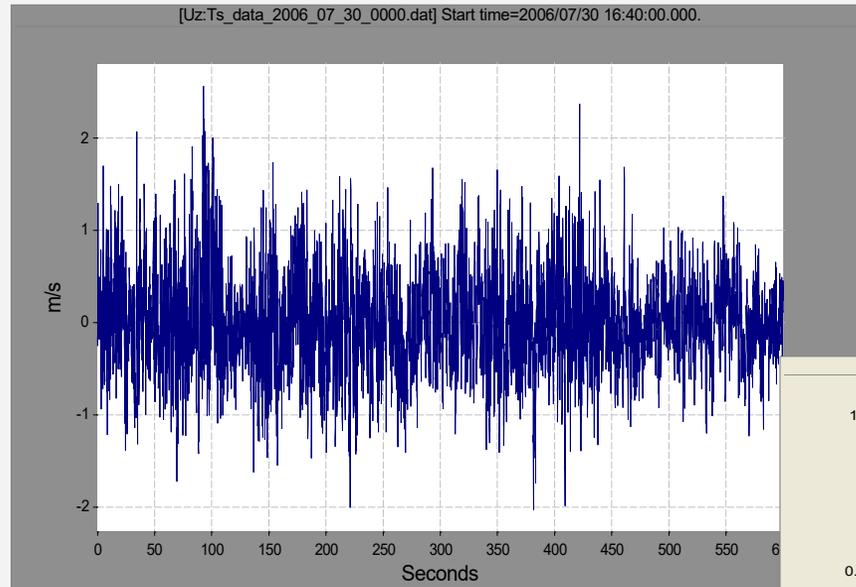
Atmospheric Boundary
Layer Flows

Their Structure and Measurement

J. C. KAIMAL
J. J. FINNIGAN

Kaimal, JC, JJ Finnigan. 1994.
Atmospheric Boundary Layer
Flows: Their Structure and
Measurement. Oxford University
Press, Oxford, 289 p.

Vertical wind (w) measured in a time domain to a frequency domain

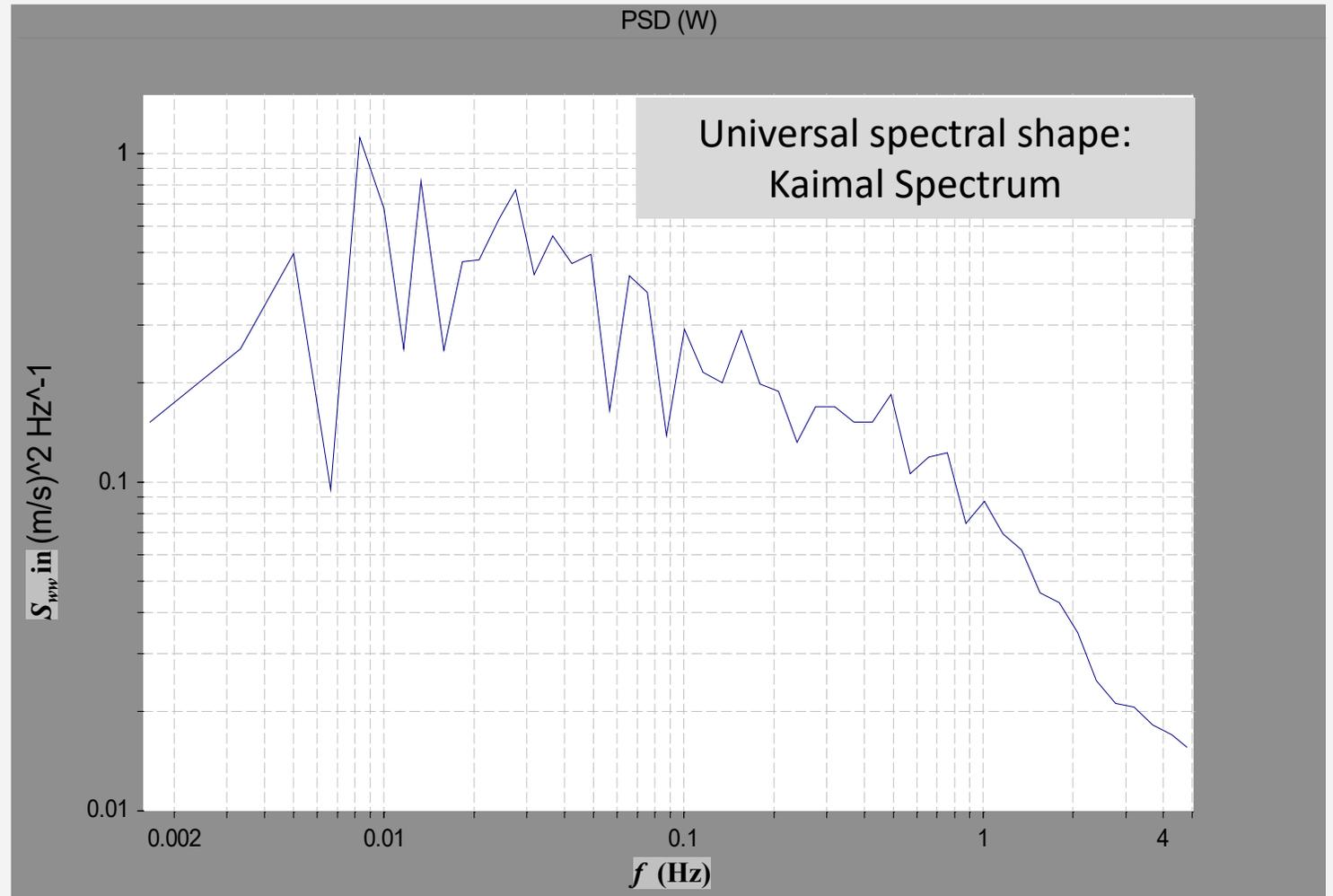


Smoothed power spectrum of vertical wind speed

$$\overline{w'w'} = \int_0^{+\infty} S_{ww}(f) df$$

Kaimal & Finnigan (1994)

Kaimal, JC, JJ Finnigan. 1994. Atmospheric Boundary Layer Flows: Their Structure and Measurement, Oxford University Press, New York, 289 p.



Relationship of cospectrum to a flux

$$\overline{\alpha'w'} = \int_{-\infty}^{+\infty} S_{\alpha w}(f) df$$

Pages 57 to 61 in Kaimal & Finnigan (1994)

α represents $u, v, T, T_s, \rho_{CO_2}, \chi_{CO_2}, \rho_{H_2O}, \chi_{H_2O}$,

OR other gas species amount.

Normalized cospectrum

$$S_{\alpha w}(f)_N = \frac{S_{\alpha w}(f)_T}{(\alpha' w')_T}$$

FREQUENCY RESPONSE CORRECTIONS FOR EDDY CORRELATION SYSTEMS

C. J. MOORE

Institute of Hydrology, Wallingford, Oxon, U.K.

(Received in final form 14 April, 1986)

Abstract. Simplified expressions describing the frequency response of eddy correlation systems due to sensor response, path-length averaging, sensor separation and signal processing are presented. A routine procedure for estimating and correcting for the frequency response loss in flux and variance measurements is discussed and illustrated by application to the Institute of Hydrology's 'Hydra' eddy correlation system.

The results show that flux loss from such a system is typically 5 to 10% for sensible and latent heat flux, but can be much larger for momentum flux and variance measurements in certain conditions.

A microcomputer program is included which, with little modification, can be used for estimating flux loss from other eddy correlation systems with different or additional sensors.

1. Introduction

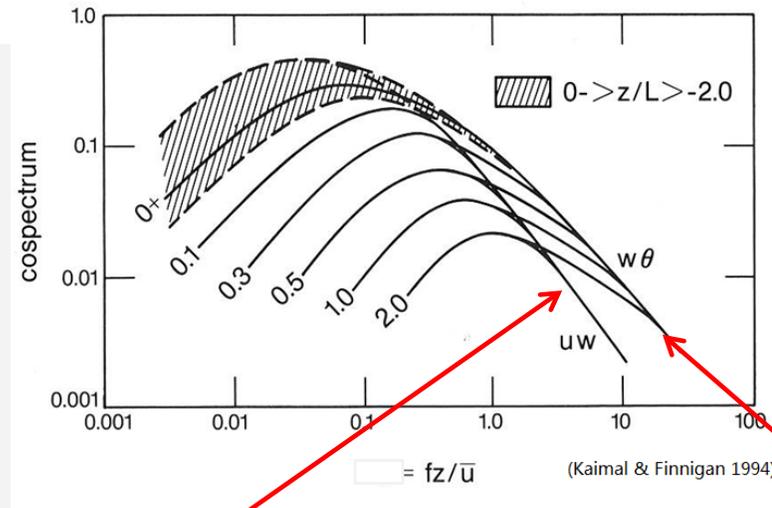
Technological advance in recent years has allowed many improvements in the design of eddy-correlation systems. This is particularly evident in the development of sonic anemometry for the measurement of wind velocity components, and several new designs have been published (e.g., Campbell and Unsworth, 1979; Larsen *et al.*, 1979; Shuttleworth *et al.*, 1982; Coppin and Taylor, 1983). Similar improvements can be found for sensors measuring atmospheric humidity (e.g., Hyson and Hicks, 1975; Raupach, 1978; Moore, 1983), carbon dioxide concentration (see Ohtaki, 1984) and other atmospheric constituents. Such development has improved the accuracy, speed of response and reliability of turbulence measurements. As a result of physical limitations in sensor size and response, experimental siting and data analysis techniques, however, these measurements will always remain frequency band-limited. Hence systematic errors leading to underestimation of turbulent fluxes and variances must be expected.

For the full potential of the eddy-correlation technique to be realized, e.g., the routine measurement of surface sensible and latent heat fluxes, it is important that the magnitude of these errors can be calculated and hence accounted for in the resulting data. Methods for estimating instrumental errors associated with particular sensors have been available for some time. Silverman (1968), Kaimal *et al.* (1968) and Horst (1973) provide path-length averaging and path-separation corrections for particular sonic anemometer arrangements frequently used in atmospheric turbulence research. Robust and inexpensive propeller anemometers have been more popular for use in eddy correlation systems measuring turbulent heat and momentum fluxes. It has been shown that measurements from these systems generally underestimate the actual fluxes by between 5 and 25% (e.g., Moore, 1976; McNeil and Shuttleworth, 1975; Spittlehouse and Black,

Moore CJ. 1986.
Frequency response corrections
for eddy correlation systems.
Bound-Layer Meteorol 37: 17-35.

Normalized Cospectra

$$\frac{z}{L} > 0$$



Vertical with horizontal wind

$$fS_{uw}(f)_N = \frac{fz/\bar{u}}{A_{uw} + B_{uw} \left(\frac{fz}{\bar{u}} \right)^{2.1}}$$

$$A_{uw} = 0.124 \left(1 + 7.9 \frac{z}{L} \right)^{0.75}$$

$$B_{uw} = 23.252 \left(1 + 7.9 \frac{z}{L} \right)^{-0.825}$$

Vertical wind with air temperature

$$fS_{Tw}(f)_N = \frac{fz/\bar{u}}{A_{Tw} + B_{Tw} \left(\frac{fz}{\bar{u}} \right)^{2.1}}$$

$$A_{Tw} = 0.284 \left(1 + 6.4 \frac{z}{L} \right)^{0.75}$$

$$B_{Tw} = 9.3447 \left(1 + 6.4 \frac{z}{L} \right)^{-0.825}$$

Normalized Cospectra

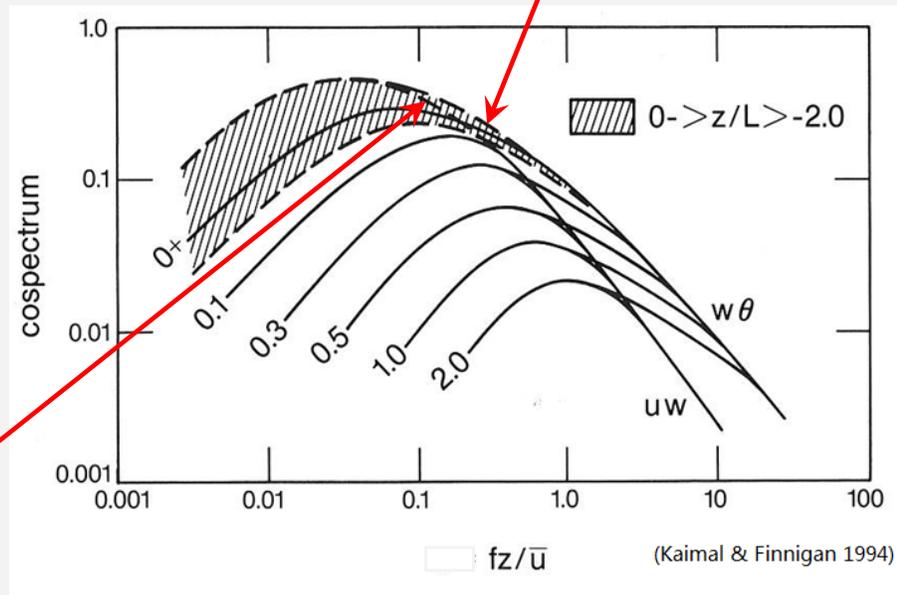
$$\frac{z}{L} \leq 0$$

Vertical with horizontal wind

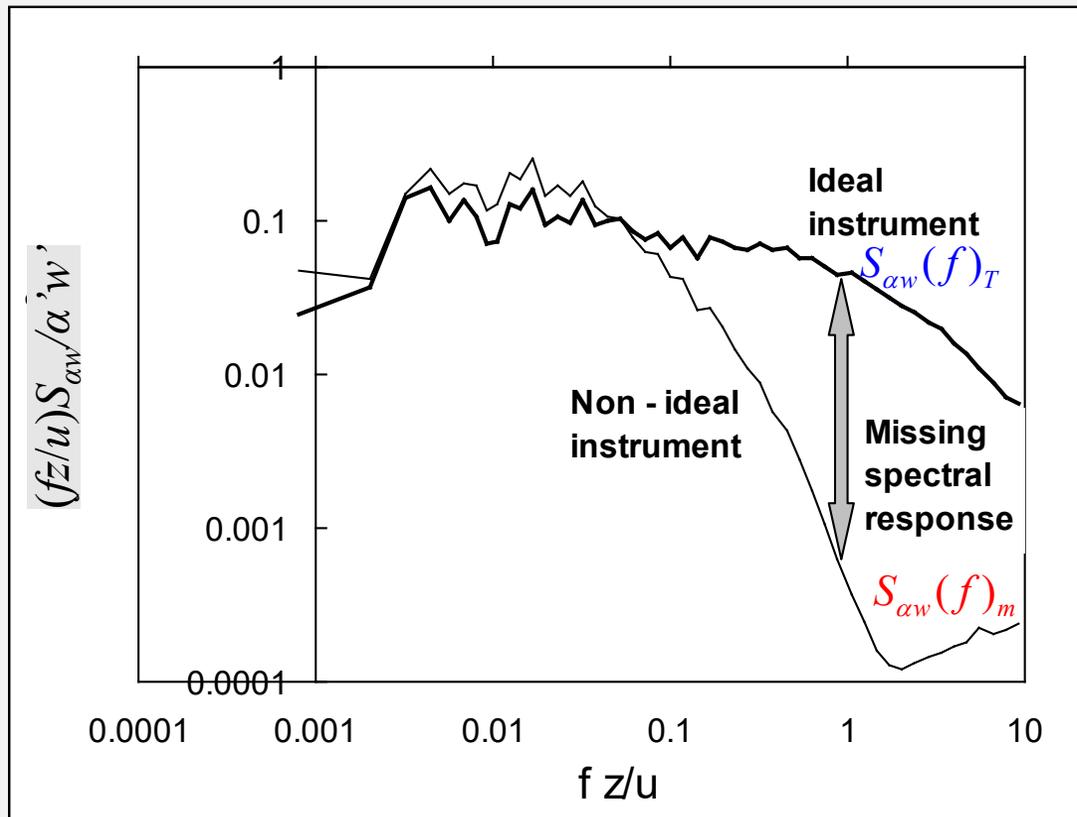
$$fS_{uw}(f)_N = \begin{cases} \frac{20.78 fz/\bar{u}}{\left(1 + \frac{31z}{\bar{u}} f\right)^{1.575}} & \frac{z}{\bar{u}} f < 0.24 \\ \frac{12.66 fz/\bar{u}}{\left(1 + \frac{9.6z}{\bar{u}} f\right)^{2.4}} & \frac{z}{\bar{u}} f \geq 0.24 \end{cases}$$

Vertical wind with air temperature

$$fS_{Tw}(f)_N = \begin{cases} \frac{12.92 fz/\bar{u}}{\left(1 + \frac{26.7z}{\bar{u}} f\right)^{1.375}} & \frac{z}{\bar{u}} f < 0.54 \\ \frac{4.378 fz/\bar{u}}{\left(1 + \frac{3.8z}{\bar{u}} f\right)^{2.4}} & \frac{z}{\bar{u}} f \geq 0.54 \end{cases}$$



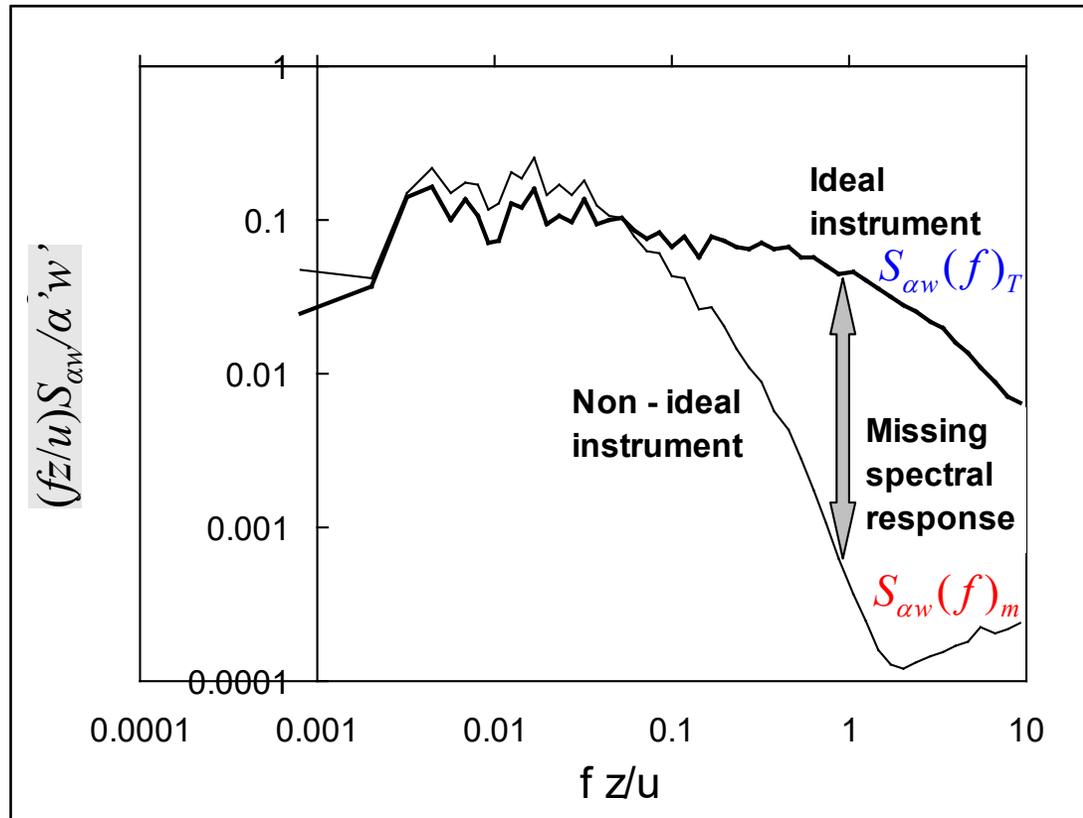
Cospectrum from an ideal vs. non-idea measurement



$$\left(\overline{\alpha' w'}\right)_T = \int_0^{\infty} S_{\alpha w}(f)_T df$$

$$\left(\overline{\alpha' w'}\right)_m = \int_0^{\infty} S_{\alpha w}(f)_m df$$

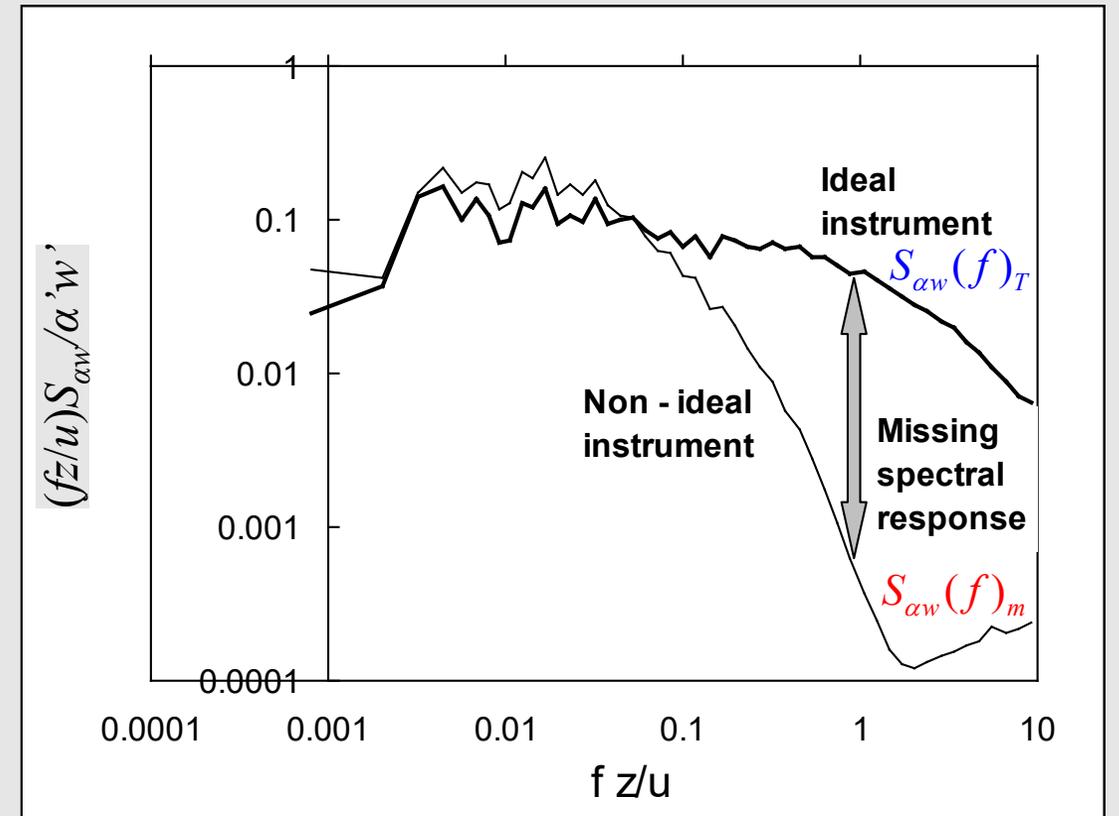
Relate “true” flux to measured flux, measured spectrum, and “true” spectrum



| | |
|---|------------------------|
| ‘true’ covariance | ‘true’ cospectrum |
| ↓ | ↓ |
| $\frac{\overline{(\alpha' w')}_T}{\overline{(\alpha' w')}_m} = \frac{\int_0^\infty S_{\alpha w}(f)_T df}{\int_0^\infty S_{\alpha w}(f)_m df}$ | |
| ↑ | ↑ |
| Measured covariance | Measured cospectrum |

Definition of Transfer Function

$$T_{\alpha w}(f) = \frac{S_{\alpha w}(f)_m}{S_{\alpha w}(f)_T} \leq 1$$



An example for the derivation of transfer function

$$\frac{dT_m(t)}{dt} = \frac{T(t) - T_m(t)}{\tau}$$

Apply Fourier Transform to both sides of this equation

$$\mathcal{F}\left(\frac{dT_m(t)}{dt}\right) = \mathcal{F}\left(\frac{T(t) - T_m(t)}{\tau}\right) \longrightarrow \begin{aligned} -i\omega T_m(\omega) &= \frac{1}{\tau} [T(\omega) - T_m(\omega)] \\ \sqrt{i} = -1 \quad \omega &= 2\pi f \end{aligned}$$

$$\frac{T_m(2\pi f)}{T(2\pi f)} = \frac{1}{1 - i\tau 2\pi f} = T_T(f)$$

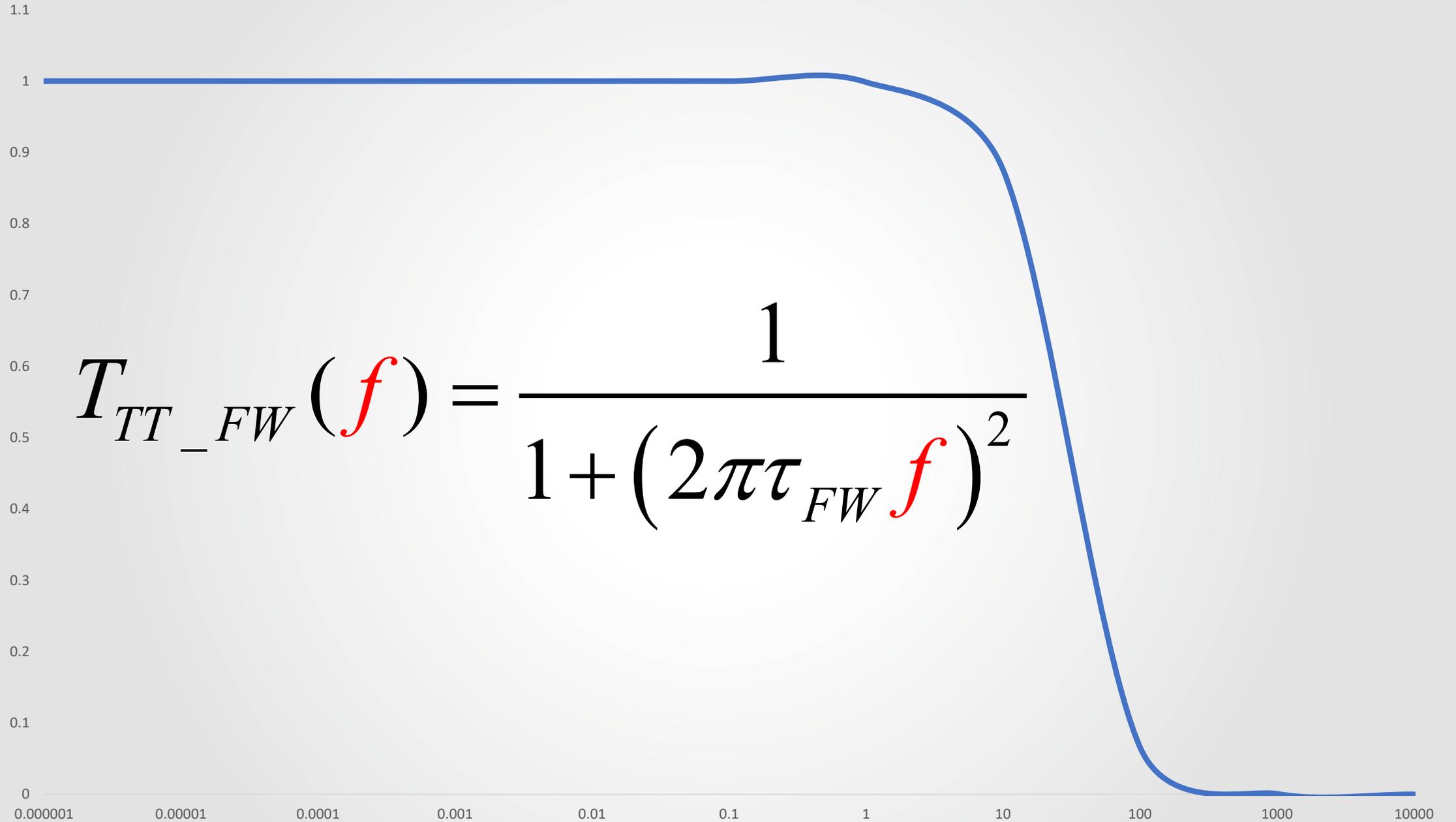
Transfer function

based on spectrum of T

$$T_T(f) = \frac{1}{1 - i2\pi\tau f}$$

based on power spectrum of T

$$T_{TT}(f) = \frac{1}{1 + (2\pi\tau f)^2}$$



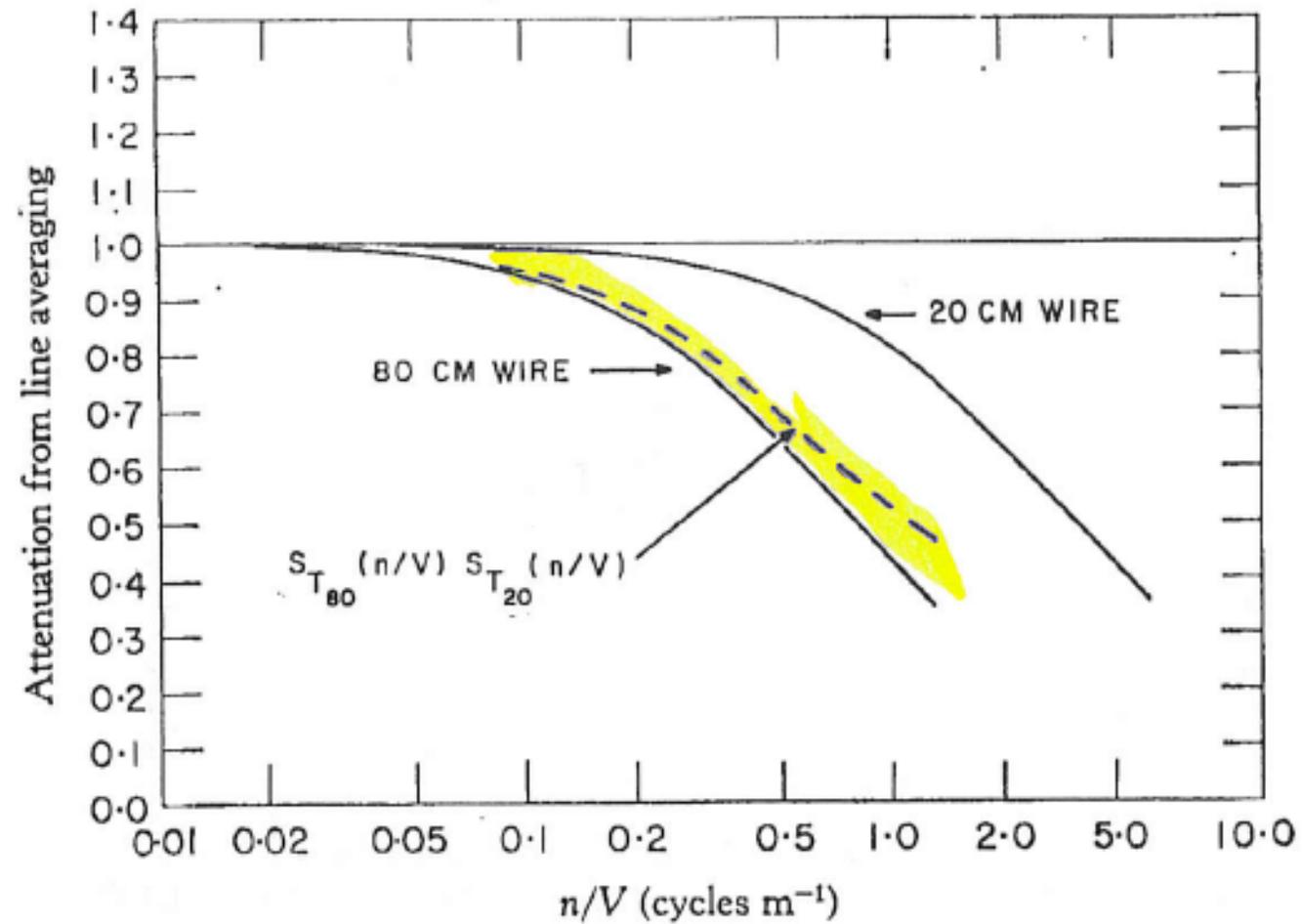


Figure 3. Theoretical attenuation curves for 20- and 80-cm line averaging. Relative attenuation between the two is indicated by dashed curves.

Line averaging

Transfer function for scalars

$$T_{ss_LA}(f) = \frac{1}{2\pi fp / u} \left(3 + \exp(-2\pi \frac{p}{u} f) - \frac{4 \left[1 - \exp(-2\pi \frac{p}{u} f) \right]}{2\pi fp / u} \right)$$

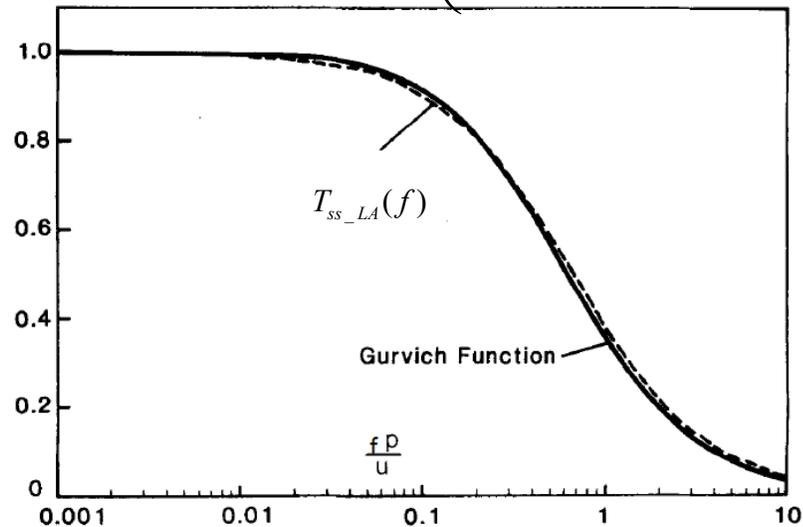
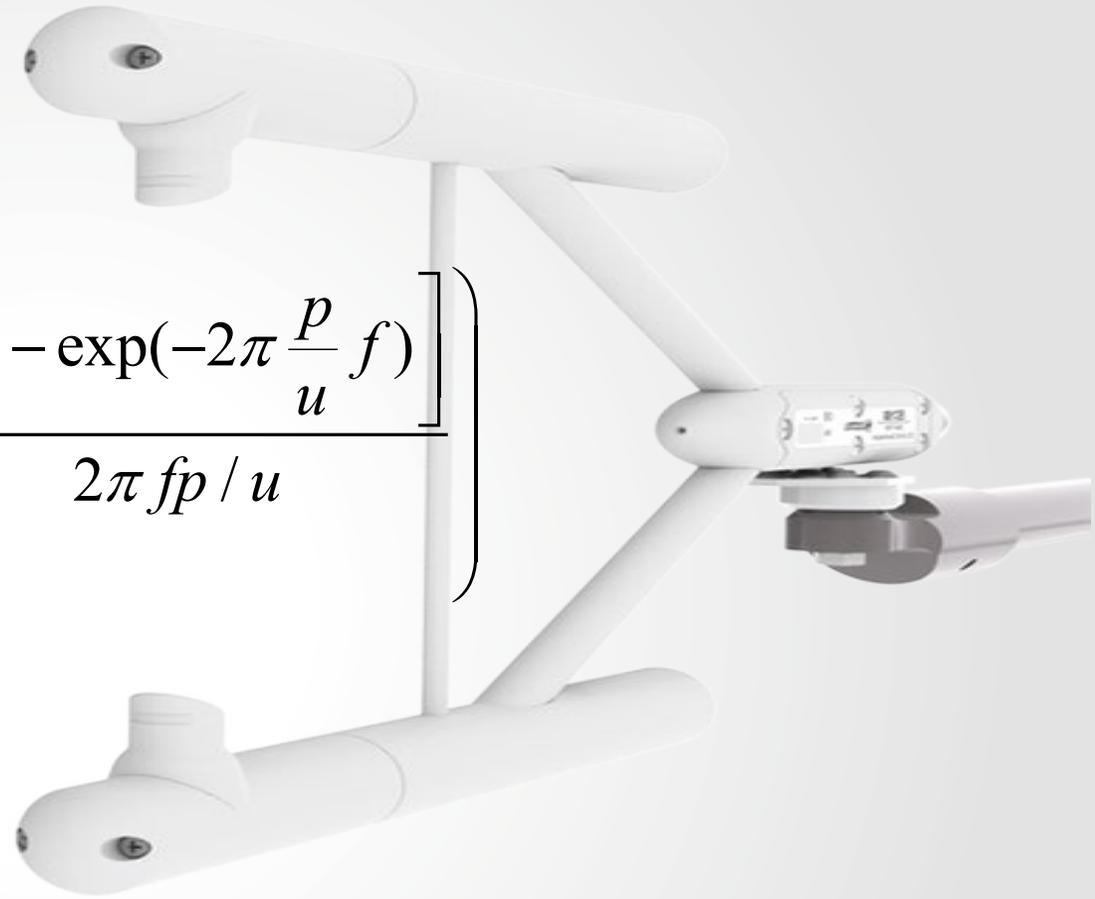


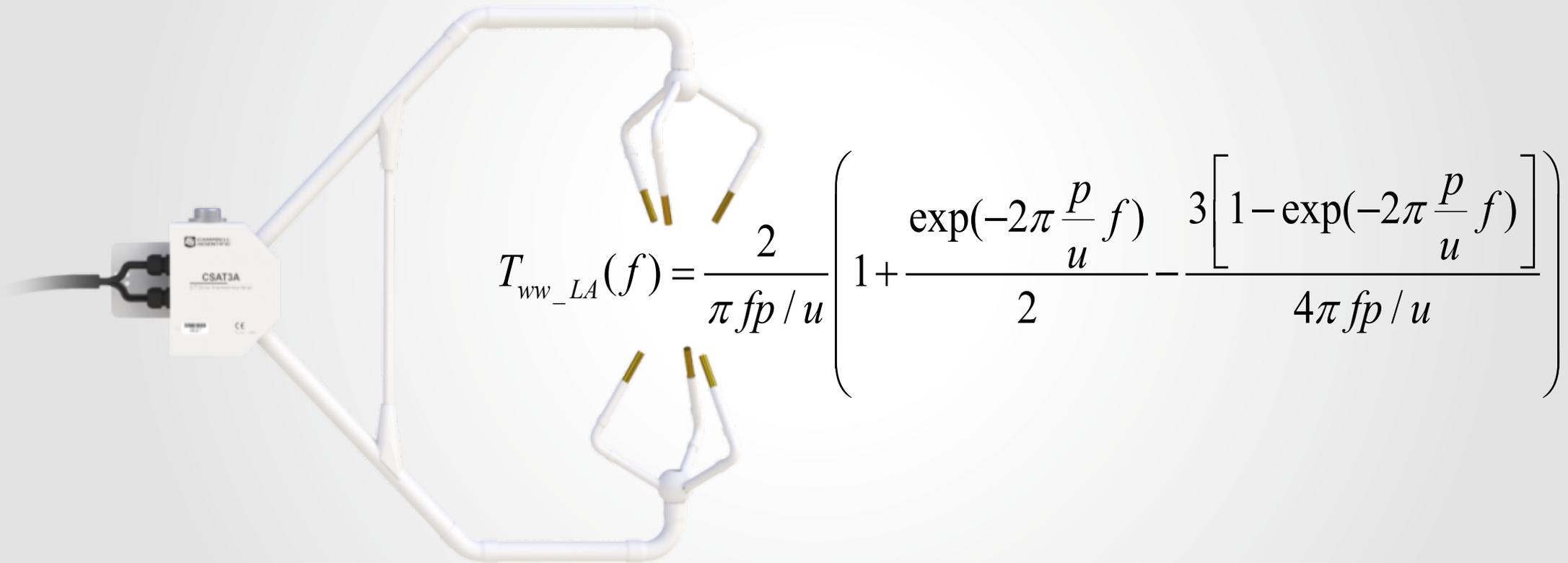
Fig. 1. Spectral transfer function, $T_{ss_LA}(f)$ associated with the measurement of a scalar quantity averaged over a finite path length, p , at right angles to the horizontal wind velocity, u , and shown as the solid curve plotted against normalized frequency fp/u . A function approximating this curve is shown as a dashed line.



Moore, CJ. 1986. Frequency response corrections for eddy correlation systems, *Boundary-Layer Meteorol*, 37: 17–35,

Line averaging

Transfer function for vertical wind



Line averaging

Transfer function for CSAT sonic temperature flux

$$T(k) = \frac{1}{\beta z w S (|k|z)^{-7/3}} \int_0^\infty \int_0^{2\pi} KA \frac{k^2 + K^2 \sin^2 \theta}{\sqrt{k^2 + K^2}} \times \frac{1}{3} \left[\sin c \left(\frac{\mathbf{k} \cdot \mathbf{I}_a}{2} \right) + \sin c \left(\frac{\mathbf{k} \cdot \mathbf{I}_b}{2} \right) + \sin c \left(\frac{\mathbf{k} \cdot \mathbf{I}_c}{2} \right) \right] d\theta dK$$

TABLE 1. Transfer functions for 1D and for 3D sonic anemometers.

| kl | 1D | 3D-30° | 3D-45° |
|------|------------|------------|------------|
| 0.01 | 1.0000 | 1.0000 | 1.0000 |
| 0.1 | 0.9998 | 0.9992 | 0.9990 |
| 0.2 | 0.9980 | 0.9976 | 0.9972 |
| 0.5 | 0.9914 | 0.9900 | 0.9886 |
| 1 | 0.9720 | 0.9670 | 0.9619 |
| 1.2 | 0.9620 | 0.9550 | 0.9479 |
| 1.4 | 0.9509 | 0.9417 | 0.9325 |
| 1.6 | 0.9390 | 0.9274 | 0.9157 |
| 1.8 | 0.9263 | 0.9122 | 0.8979 |
| 2 | 0.9130 | 0.8962 | 0.8792 |
| 2.2 | 0.8993 | 0.8797 | 0.8597 |
| 2.4 | 0.8851 | 0.8626 | 0.8397 |
| 2.6 | 0.8707 | 0.8452 | 0.8192 |
| 2.8 | 0.8560 | 0.8274 | 0.7983 |
| 3 | 0.8412 | 0.8096 | 0.7773 |
| 4 | 0.7668 | 0.7201 | 0.6720 |
| 5 | 0.6954 | 0.6353 | 0.5729 |
| 6 | 0.6296 | 0.5588 | 0.4855 |
| 7 | 0.5706 | 0.4922 | 0.4117 |
| 8 | 0.5185 | 0.4355 | 0.3515 |
| 9 | 0.4728 | 0.3879 | 0.3038 |
| 10 | 0.4330 | 0.3481 | 0.2664 |
| 14 | 0.3181 | 0.2445 | 0.1822 |
| 20 | 0.2241 | 0.1700 | 0.1288 |
| 30 | 0.1495 | 0.1134 | 0.8580E-01 |
| 40 | 0.1121 | 0.8503E-01 | 0.6436E-01 |
| 50 | 0.8968E-01 | 0.6802E-01 | 0.5149E-01 |
| 60 | 0.7473E-01 | 0.5668E-01 | 0.4291E-01 |
| 70 | 0.6406E-01 | 0.4859E-01 | 0.3678E-01 |
| 80 | 0.5605E-01 | 0.4251E-01 | 0.3218E-01 |
| 90 | 0.4982E-01 | 0.3779E-01 | 0.2860E-01 |
| 100 | 0.4484E-01 | 0.3401E-01 | 0.2574E-01 |

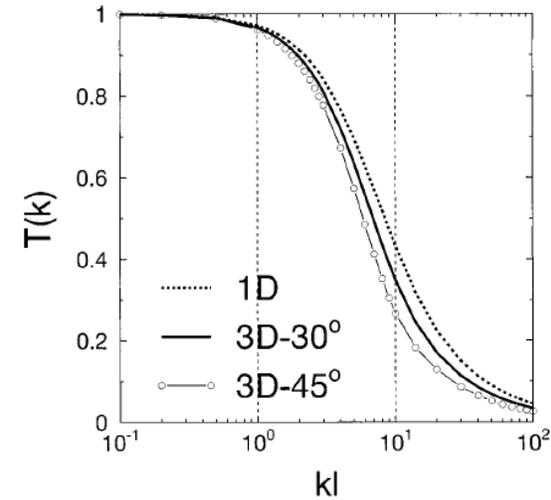


FIG. 3. Transfer functions for sonic anemometers: the dashed line gives the 1D case, the continuous line is for the 3D Campbell-Solent type of sonic and the circles represent a configuration where the acoustic paths make an angle of 45° with the vertical axis. The quantity on the horizontal axis is wavenumber $k = 2\pi/\lambda$ of a flux contribution made dimensionless with acoustic pathlength l .

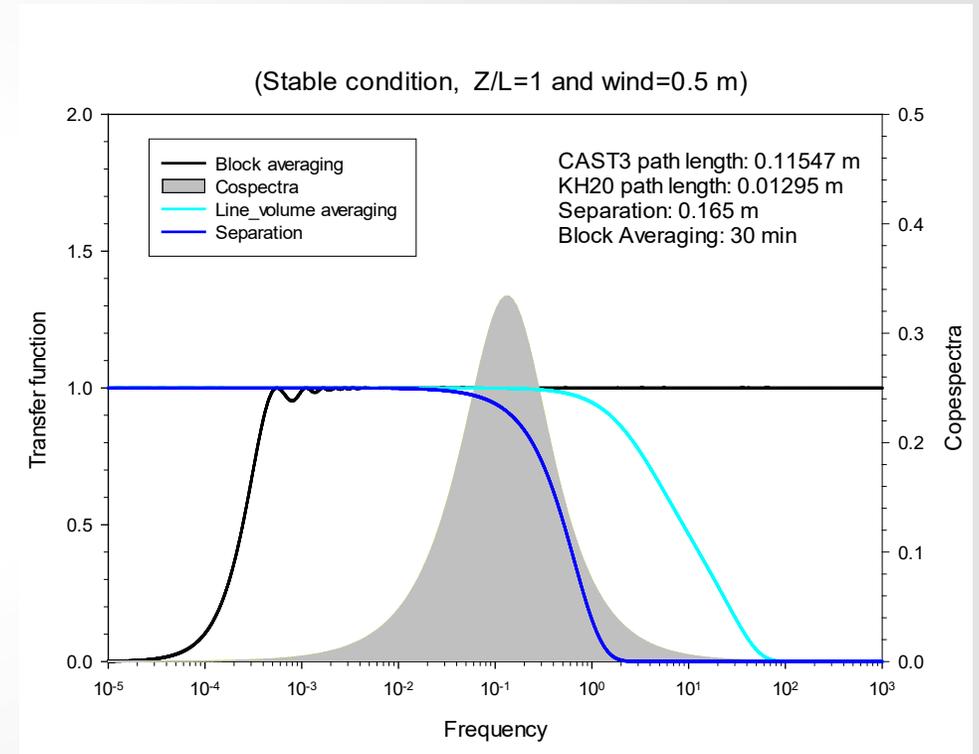
van Dijk (2002)

Separation between sensors

Transfer functions for vertical wind with scalar

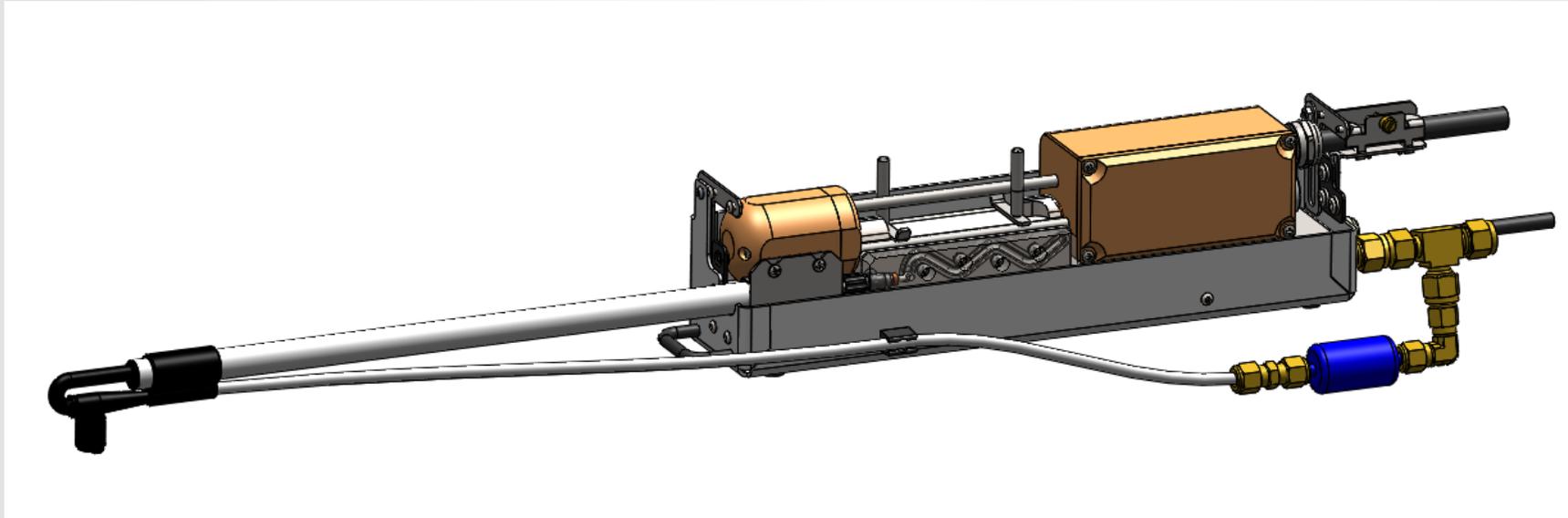
$$T_{w_s}(f) = \exp \left[-9.9 \left(\frac{f \times d}{\bar{u}} \right)^{1.5} \right]$$

Alternatively, lag maximization



Response delay

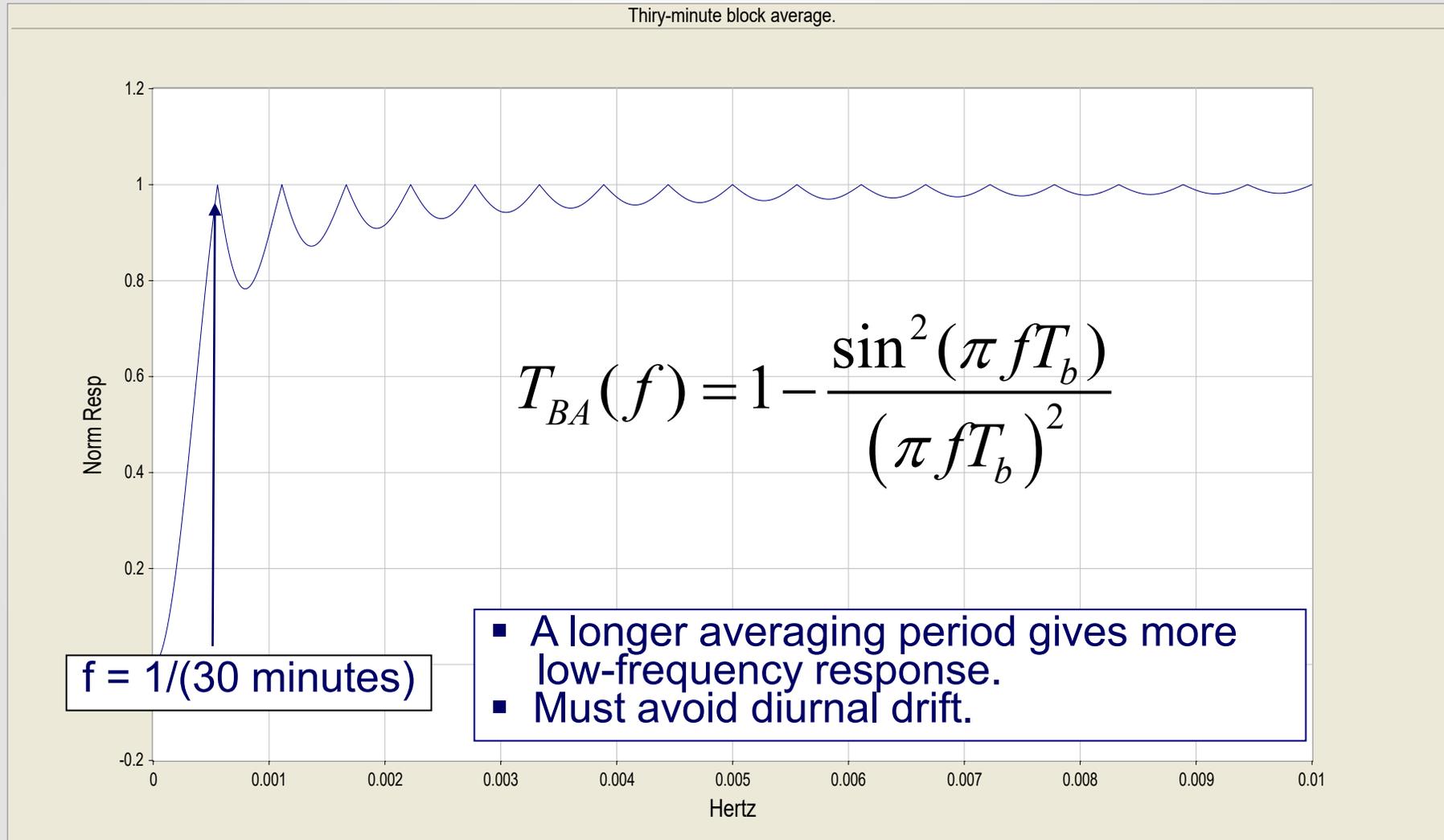
Transfer functions for closed-path sampling cell



$$T_{CO_2CO_2}(f) = \frac{1}{1 + \left(2\pi \frac{f}{4.6}\right)^2}$$

Transfer function

block averaging



Derivation of measured cospectrum

$$T_{\alpha w}(f) = \frac{S_{\alpha w}(f)_m}{S_{\alpha w}(f)_T} \leq 1$$

$$S_{\alpha w}(f)_m = T_{\alpha w}(f) S_{\alpha w}(f)_T$$

$$S_{\alpha w}(f)_m = T_{\alpha w}(f) S_{\alpha w}(f)_T$$

'true'
covariance



$$\frac{\overline{(\alpha' w')}_T}{\overline{(\alpha' w')}_m} = \frac{\int_0^\infty S_{\alpha w}(f)_T df}{\int_0^\infty S_{\alpha w}(f)_m df}$$

Measured
covariance



'true'
cospectrum



$$\frac{\overline{(\alpha' w')}_T}{\overline{(\alpha' w')}_m} = \frac{\int_0^\infty S_{\alpha w}(f)_T df}{\int_0^\infty T_{\alpha w}(f) S_{\alpha w}(f)_T df}$$

Measured
cospectrum



Transfer
function



$$\frac{\left(\overline{\alpha'w'}\right)_T}{\left(\overline{\alpha'w'}\right)_m} = \frac{\int_0^\infty S_{\alpha w}(f)_T df}{\int_0^\infty T_{\alpha w}(f) S_{\alpha w}(f)_T df} = \frac{\int_0^\infty \left[\left(\overline{\alpha'w'}\right)_T S_{\alpha w}(f)_N \right] df}{\int_0^\infty T_{\alpha w}(f) \left[\left(\overline{\alpha'w'}\right)_T S_{\alpha w}(f)_N \right] df}$$

**'true'
covariance**



$$\left(\overline{\alpha'w'}\right)_T = \left(\overline{\alpha'w'}\right)_m \frac{\int_0^\infty S_{\alpha w}(f)_N df}{\int_0^\infty T_{\alpha w}(f) S_{\alpha w}(f)_N df}$$



**Measured
covariance**



**Transfer
function**

**'true'
covariance**

$$\left(\overline{\alpha'w'}\right)_T = \left(\overline{\alpha'w'}\right)_m$$

**Measured
covariance**

$$\frac{\int_0^\infty S_{\alpha w}(f)_N df}{\int_0^\infty T_{\alpha w}(f) S_{\alpha w}(f)_N df}$$

**Transfer
function**

Frequency
correction factor
 f_c

$$f_c = \frac{\int_0^\infty S_{\alpha w}(f)_N df}{\int_0^\infty T_{\alpha w}(f) S_{\alpha w}(f)_N df} \geq 1$$

$$\left(\overline{\alpha'w'}\right)_T = f_c \left(\overline{\alpha'w'}\right)_m$$

Frequency correction factor

$$f_c = \frac{\int_0^{\infty} S_{\alpha w}(f)_N df}{\int_0^{\infty} T_{\alpha w}(f) S_{\alpha w}(f)_N df} \geq 1$$


**A product of
all transfer functions
for related sensors.**

$$\left(\overline{\alpha' w'} \right)_T = f_c \left(\overline{\alpha' w'} \right)_m$$

The composite Simpson's rule writes the integration

$$\int_{10^{-6}}^{10000} f S_{\alpha w}(f)_N d(\ln f) \approx \frac{\ln \Delta f}{3} \left\{ \left[10^{-6} S_{\alpha w}(10^{-6})_N \right] + 4 \sum_{k=1}^{50} \left[10^{-6} \Delta f^{2k-1} S_{\alpha w}(10^{-6} \Delta f^{2k-1})_N \right] \right. \\ \left. + 2 \sum_{k=1}^{49} \left[10^{-6} \Delta f^{2k} S_{\alpha w}(10^{-6} \Delta f^{2k})_N \right] + \left[10^4 S_{\alpha w}(10^4)_N \right] \right\}$$

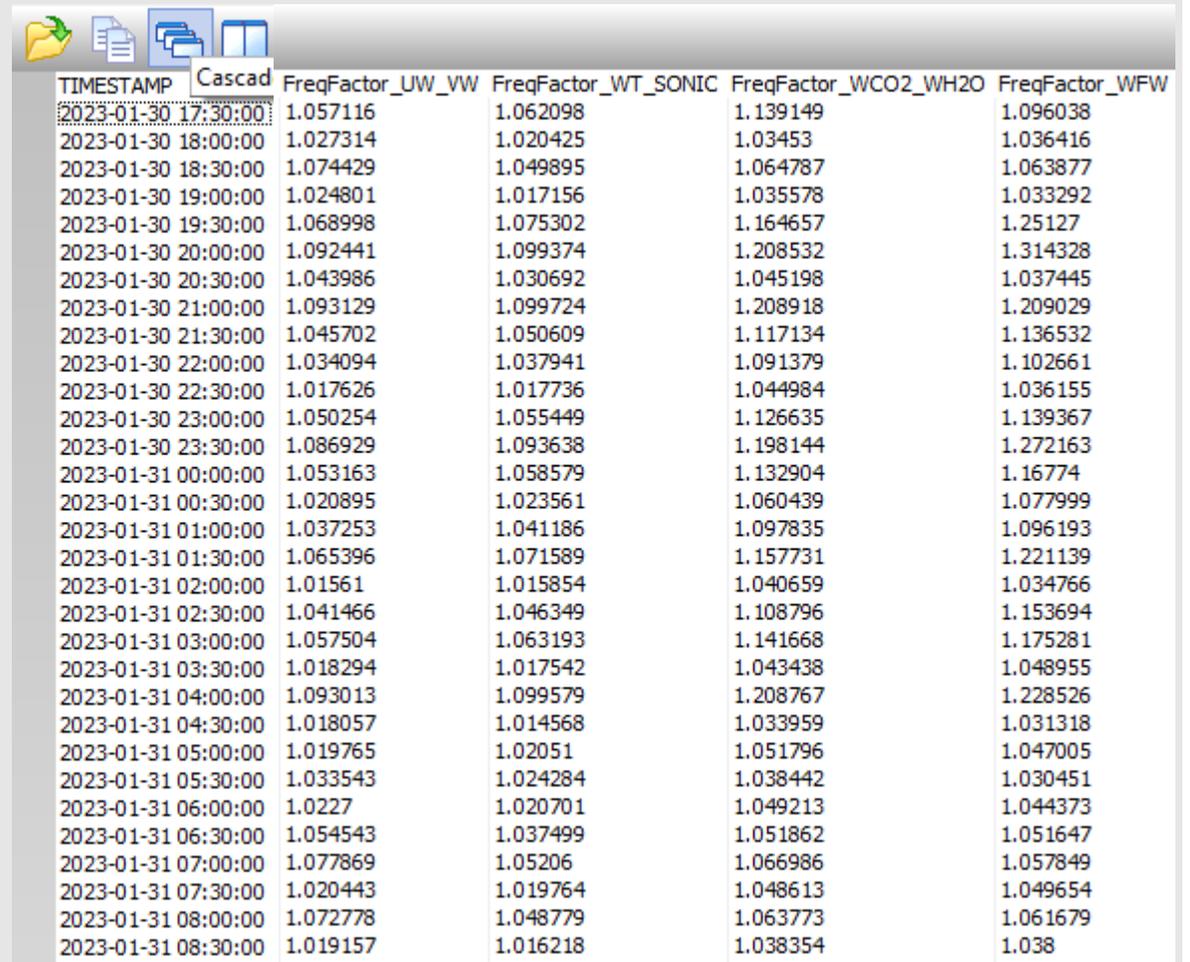
$$f_c = \frac{\int_{10^{-6}}^{10000} f S_{\alpha w}(f)_N d \ln f}{\int_{10^{-6}}^{10000} T_{\alpha w}(f) [f S_{\alpha w}(f)_N] d \ln f}$$

Burden and Faires (1993)

$$\int_{10^{-6}}^{10000} T_{\alpha w}(f) [f S_{\alpha w}(f)_N] d(\ln f) \approx \frac{\ln \Delta f}{3} \left\{ \left[T_{\alpha w}(10^{-6}) 10^{-6} S_{\alpha w}(10^{-6})_N \right] + 4 \sum_{k=1}^{50} T_{\alpha w}(10^{-6} \Delta f^{2k-1}) \left[10^{-6} \Delta f^{2k-1} S_{\alpha w}(10^{-6} \Delta f^{2k-1})_N \right] \right. \\ \left. + 2 \sum_{k=1}^{49} T_{\alpha w}(10^{-6} \Delta f^{2k}) \left[10^{-6} \Delta f^{2k} S_{\alpha w}(10^{-6} \Delta f^{2k})_N \right] + T_{\alpha w}(10^4) \left[10^4 S_{\alpha w}(10^4)_N \right] \right\}$$

f_c

Correction factors in Flux_notes
file for different covariance
(i.e., flux)



The screenshot shows a spreadsheet with the following columns: TIMESTAMP, Cascad, FreqFactor_UW_VW, FreqFactor_WT_SONIC, FreqFactor_WCO2_WH2O, and FreqFactor_WFW. The data spans from 2023-01-30 17:30:00 to 2023-01-31 08:30:00 in 30-minute intervals.

| TIMESTAMP | Cascad | FreqFactor_UW_VW | FreqFactor_WT_SONIC | FreqFactor_WCO2_WH2O | FreqFactor_WFW |
|---------------------|--------|------------------|---------------------|----------------------|----------------|
| 2023-01-30 17:30:00 | | 1.057116 | 1.062098 | 1.139149 | 1.096038 |
| 2023-01-30 18:00:00 | | 1.027314 | 1.020425 | 1.03453 | 1.036416 |
| 2023-01-30 18:30:00 | | 1.074429 | 1.049895 | 1.064787 | 1.063877 |
| 2023-01-30 19:00:00 | | 1.024801 | 1.017156 | 1.035578 | 1.033292 |
| 2023-01-30 19:30:00 | | 1.068998 | 1.075302 | 1.164657 | 1.25127 |
| 2023-01-30 20:00:00 | | 1.092441 | 1.099374 | 1.208532 | 1.314328 |
| 2023-01-30 20:30:00 | | 1.043986 | 1.030692 | 1.045198 | 1.037445 |
| 2023-01-30 21:00:00 | | 1.093129 | 1.099724 | 1.208918 | 1.209029 |
| 2023-01-30 21:30:00 | | 1.045702 | 1.050609 | 1.117134 | 1.136532 |
| 2023-01-30 22:00:00 | | 1.034094 | 1.037941 | 1.091379 | 1.102661 |
| 2023-01-30 22:30:00 | | 1.017626 | 1.017736 | 1.044984 | 1.036155 |
| 2023-01-30 23:00:00 | | 1.050254 | 1.055449 | 1.126635 | 1.139367 |
| 2023-01-30 23:30:00 | | 1.086929 | 1.093638 | 1.198144 | 1.272163 |
| 2023-01-31 00:00:00 | | 1.053163 | 1.058579 | 1.132904 | 1.16774 |
| 2023-01-31 00:30:00 | | 1.020895 | 1.023561 | 1.060439 | 1.077999 |
| 2023-01-31 01:00:00 | | 1.037253 | 1.041186 | 1.097835 | 1.096193 |
| 2023-01-31 01:30:00 | | 1.065396 | 1.071589 | 1.157731 | 1.221139 |
| 2023-01-31 02:00:00 | | 1.01561 | 1.015854 | 1.040659 | 1.034766 |
| 2023-01-31 02:30:00 | | 1.041466 | 1.046349 | 1.108796 | 1.153694 |
| 2023-01-31 03:00:00 | | 1.057504 | 1.063193 | 1.141668 | 1.175281 |
| 2023-01-31 03:30:00 | | 1.018294 | 1.017542 | 1.043438 | 1.048955 |
| 2023-01-31 04:00:00 | | 1.093013 | 1.099579 | 1.208767 | 1.228526 |
| 2023-01-31 04:30:00 | | 1.018057 | 1.014568 | 1.033959 | 1.031318 |
| 2023-01-31 05:00:00 | | 1.019765 | 1.02051 | 1.051796 | 1.047005 |
| 2023-01-31 05:30:00 | | 1.033543 | 1.024284 | 1.038442 | 1.030451 |
| 2023-01-31 06:00:00 | | 1.0227 | 1.020701 | 1.049213 | 1.044373 |
| 2023-01-31 06:30:00 | | 1.054543 | 1.037499 | 1.051862 | 1.051647 |
| 2023-01-31 07:00:00 | | 1.077869 | 1.05206 | 1.066986 | 1.057849 |
| 2023-01-31 07:30:00 | | 1.020443 | 1.019764 | 1.048613 | 1.049654 |
| 2023-01-31 08:00:00 | | 1.072778 | 1.048779 | 1.063773 | 1.061679 |
| 2023-01-31 08:30:00 | | 1.019157 | 1.016218 | 1.038354 | 1.038 |

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Questions?

